

Measurement of $(n,xn\gamma)$ reactions of interest for the new nuclear reactors

Jean Claude Thiry

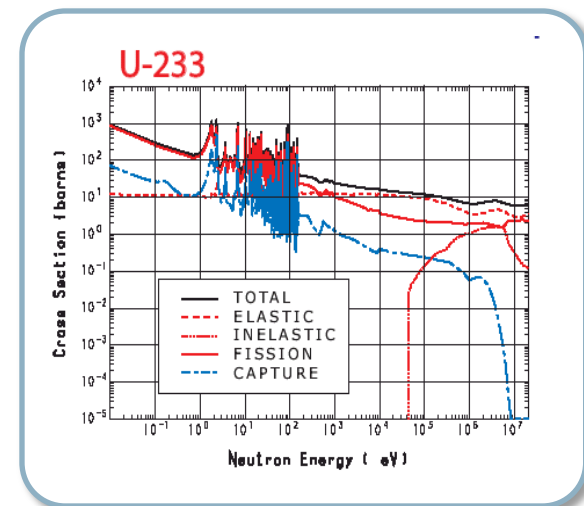
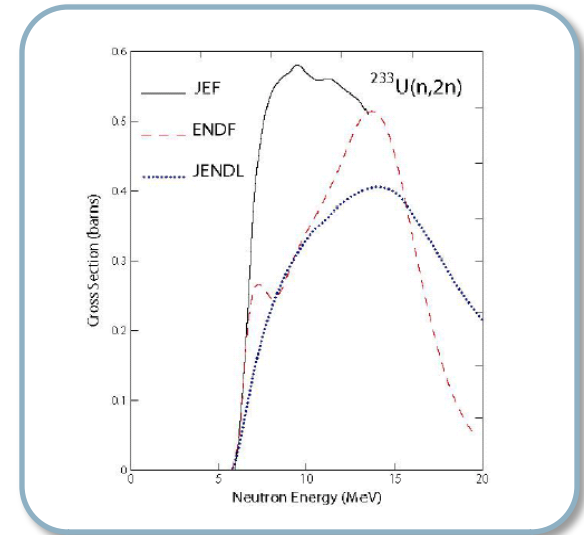


In2p3



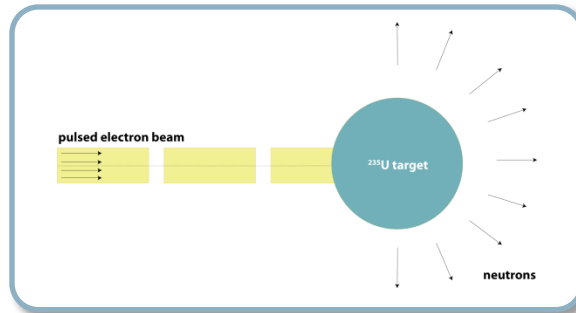
Motivation

- New reactor generation based on fast neutron reactions:
 - neutron induced reactions in these regions not known or not precisely
- Importance of (n,xn) reactions
 - neutron slow down
 - neutron multiplication
 - production of other nuclei
- Measurements on ^{206}Pb , ^{207}Pb , ^{208}Pb , ^{232}Th have been made. Current work: ^{235}U , natW



Experimental setup

- Experiments are performed using a white pulsed neutron beam from Gelina (IRMM, Geel, Belgium)
- A ^{235}U target is bombarded by a pulsed electron beam: γ radiation is created by *Bremsstrahlung*, neutrons are created by *photofission*.



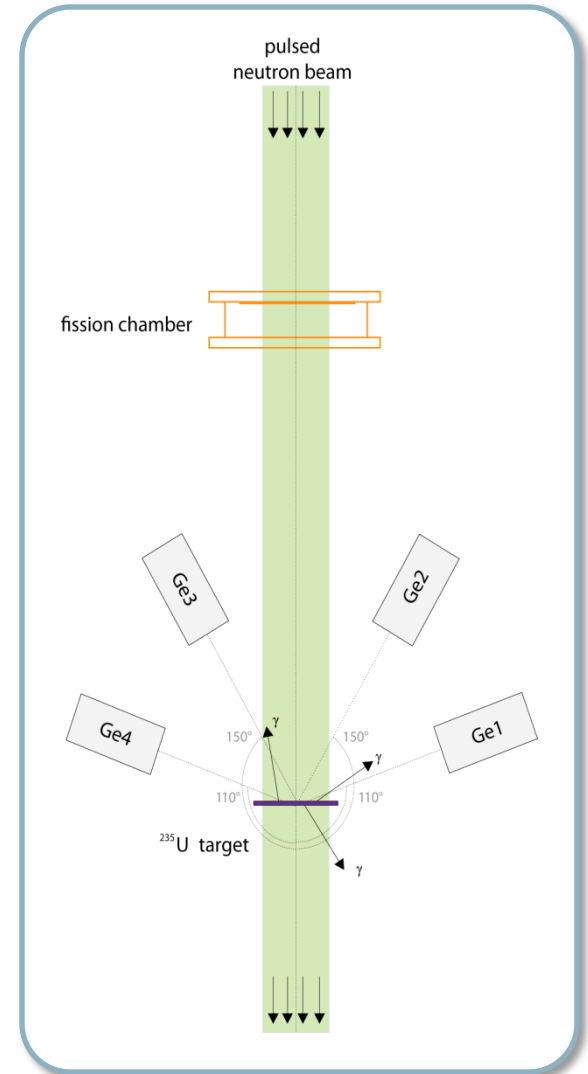
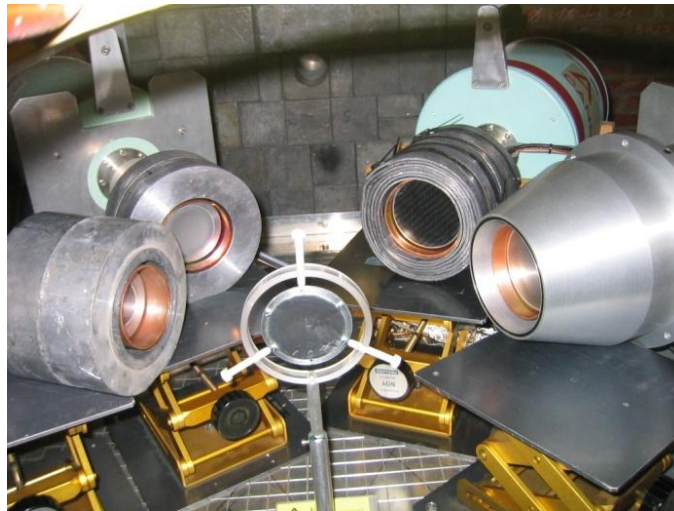
- Neutron energy is determined using the TOF technique

$$E_n = (\gamma - 1)m_n c^2 \quad , \quad \gamma = \sqrt{\frac{1}{1 - \frac{v^2}{c^2}}} \quad , \quad v = \frac{d}{t}$$



Experimental setup

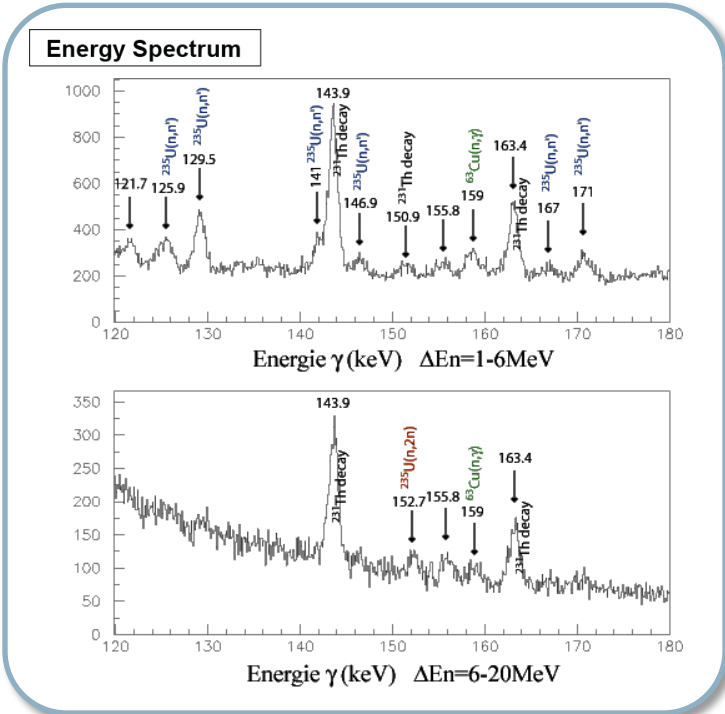
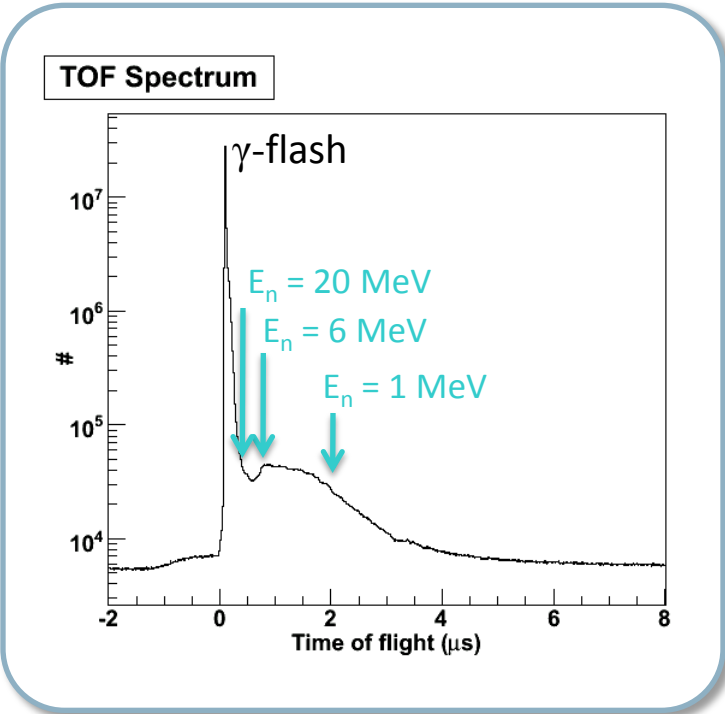
- Experiments are performed at FP16/30m
- Method: Prompt γ ray spectroscopy
 - Ex: $^{235}\text{U}(n,n')^{235}\text{U}^*$
 $^{235}\text{U}(n,2n)^{234}\text{U}^*$



- DAQ: Digital TNT2 cards developed at IPHC

HPGe detectors

- Time of flight and energy spectra



In2p3

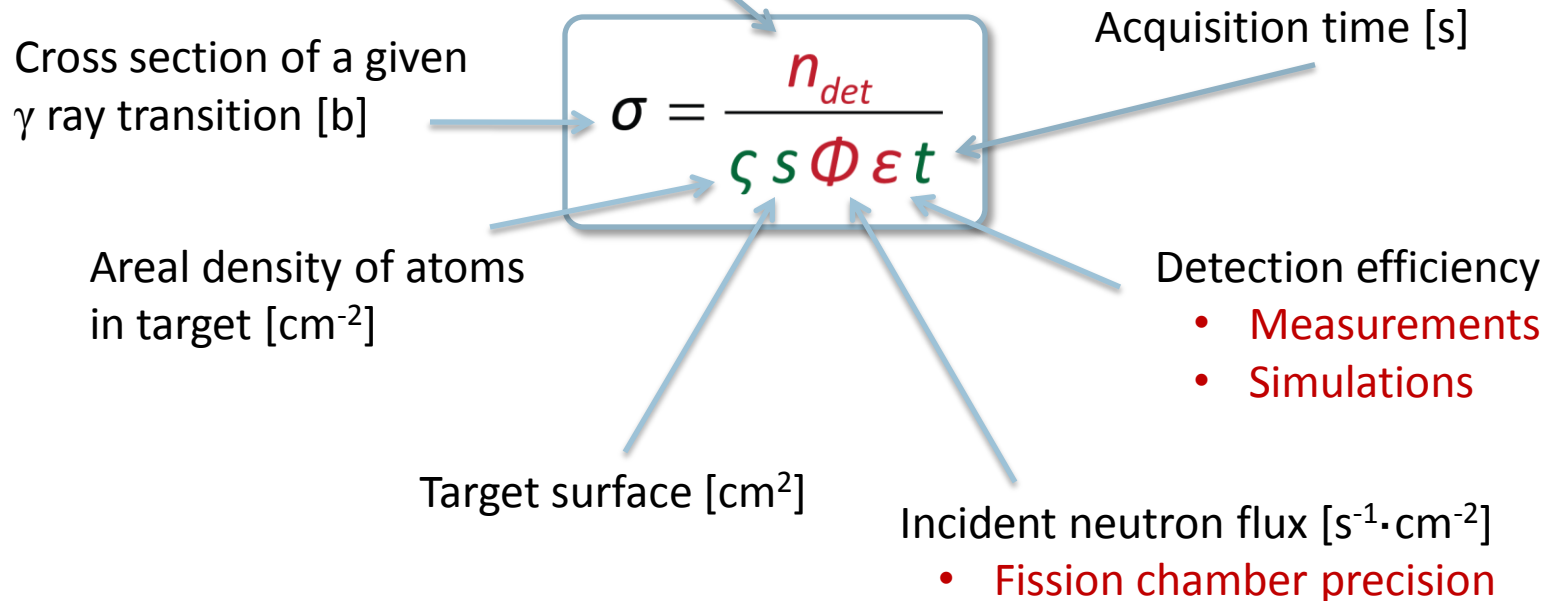


Cross sections

- The cross section of a γ transition can be expressed as:

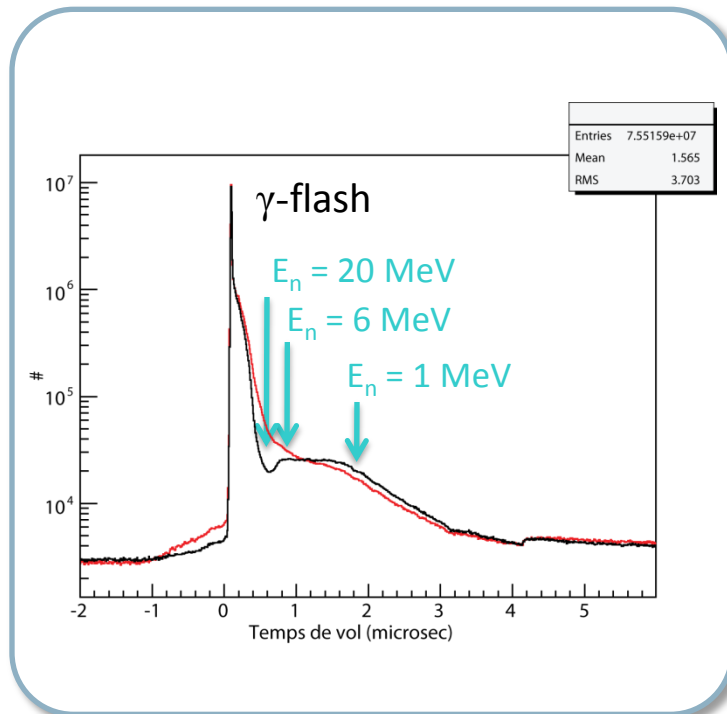
Detected hits in a given ray:

- Good statistics
- Low background



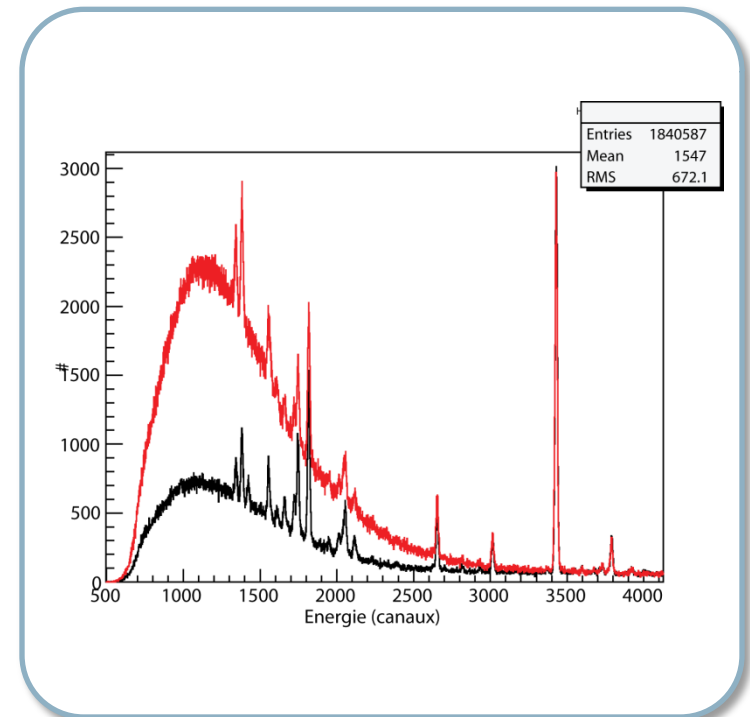
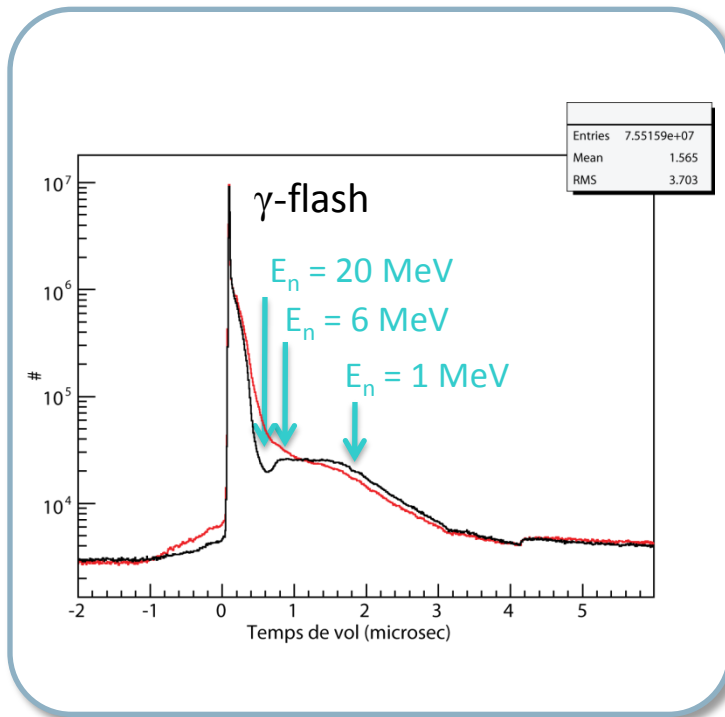
Background reduction in Ge spectra

- Time of flight and energy spectra



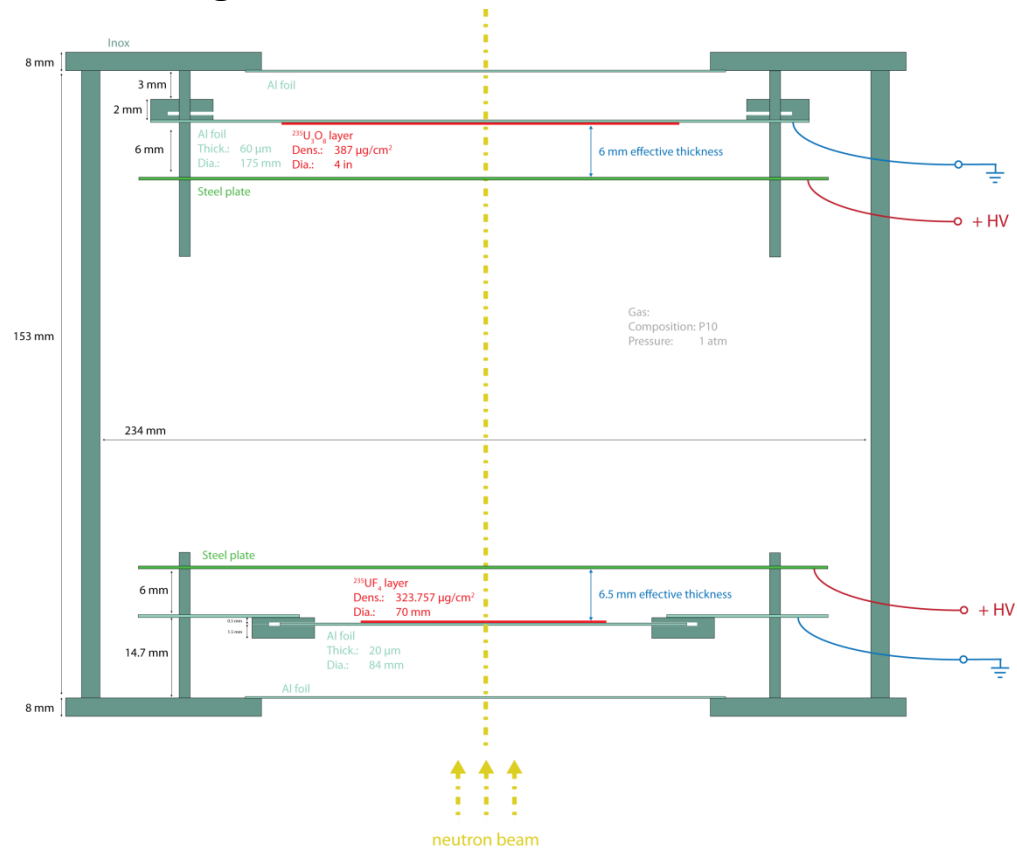
Background reduction in Ge spectra

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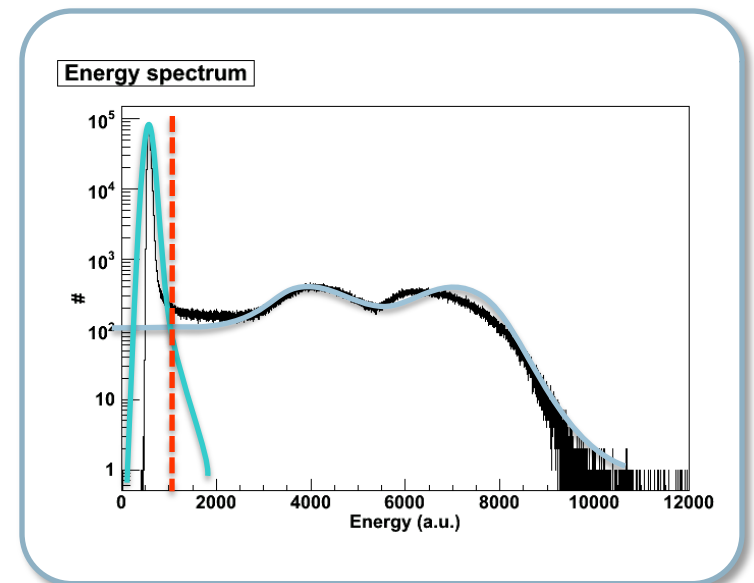
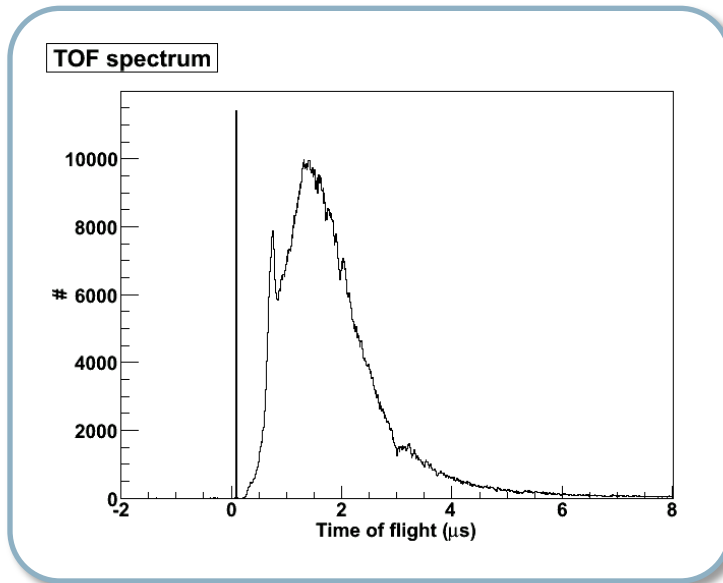
Flux monitoring

- To get accurate results it is of utmost importance to know precisely the incident neutron flux
 - Flux monitored using a fission chamber



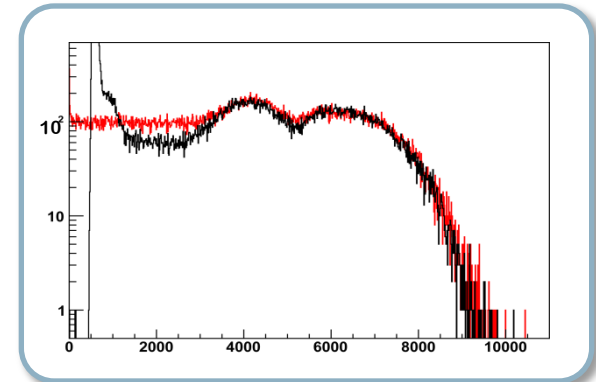
Fission chamber

- Time of flight and energy spectra

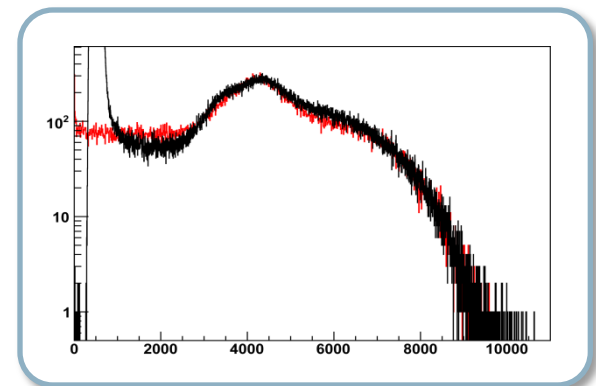


Simulating the fission chamber

- On one hand some of the created fission products don't have sufficient energy to get out of the $^{235}\text{UF}_4$ foil
 - They don't generate a detectable pulse
- On the other hand, weak signals are lost by performing a straight cut at a certain energy level
- To take these facts into account, simulations have been made using the Geant4 code
- The energies of the fission products are too weak (below the Bragg peak) for correct Geant4 simulations
- Calibration measurements will have to be performed: current estimated efficiency of $\approx 85\%$



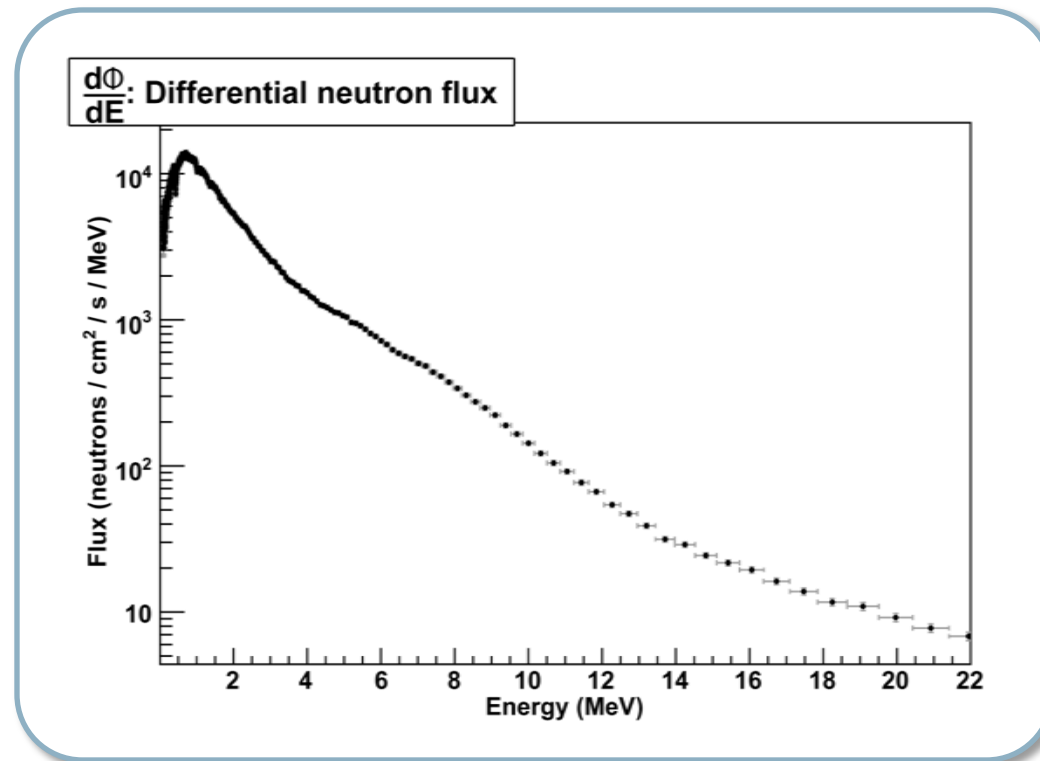
Simulated energy spectrum for a FC thickness of 8 mm



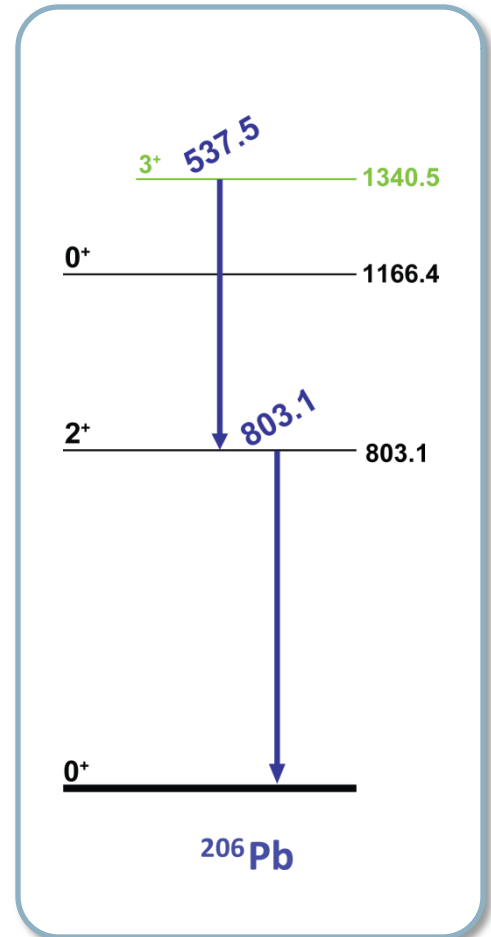
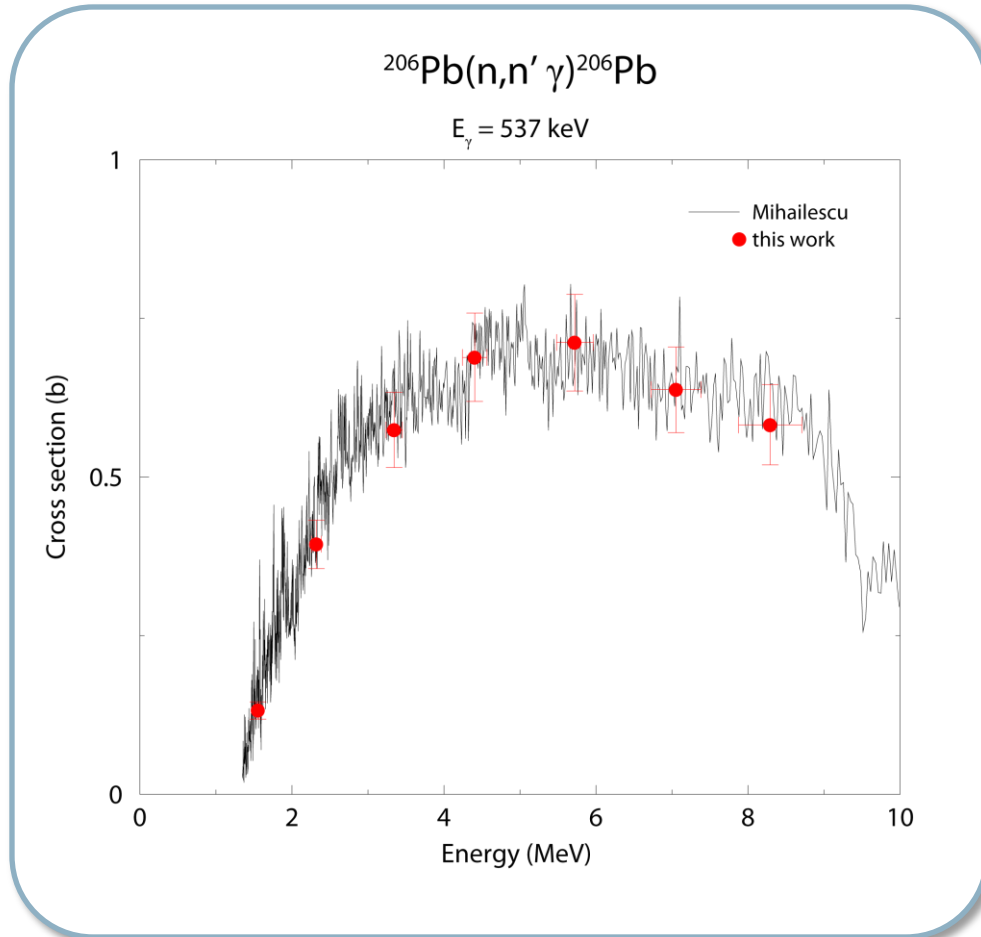
Simulated energy spectrum for a FC thickness of 5.5 mm

Neutron flux at Gelina

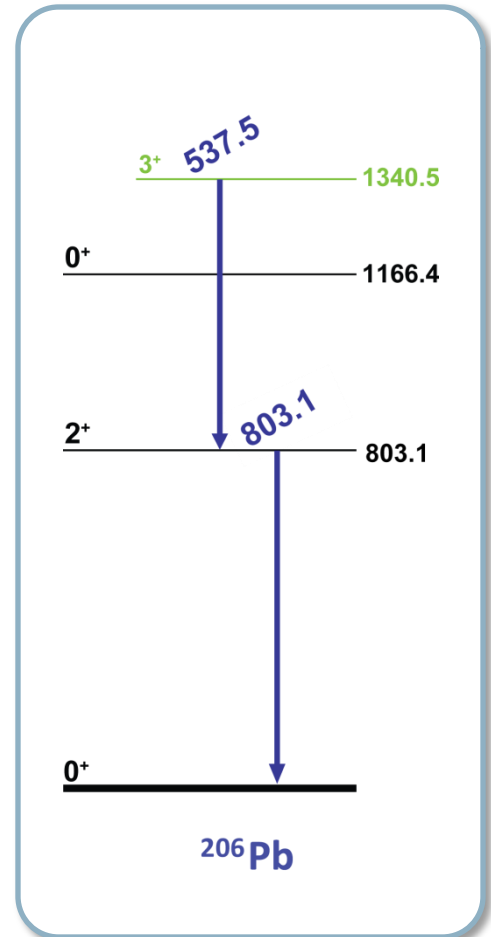
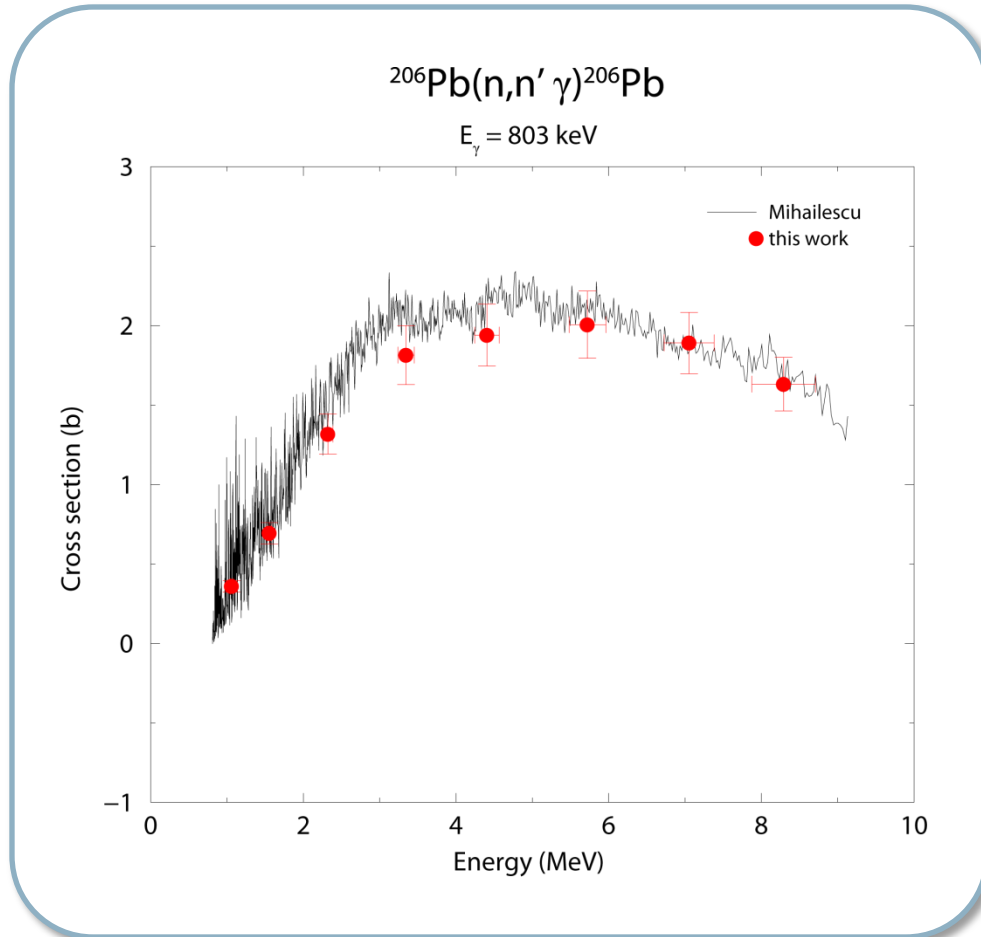
- The given neutron flux is highest at 1 MeV
- Uncertainties currently range from 8 to 10 %



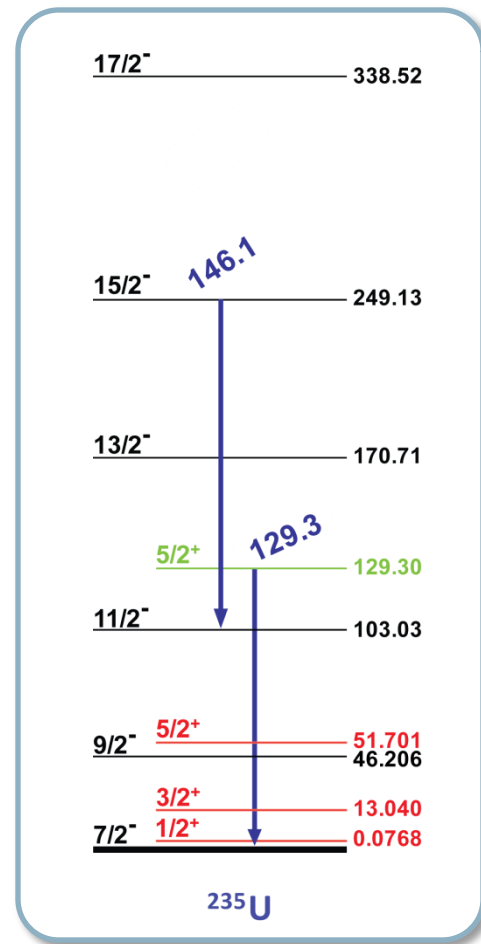
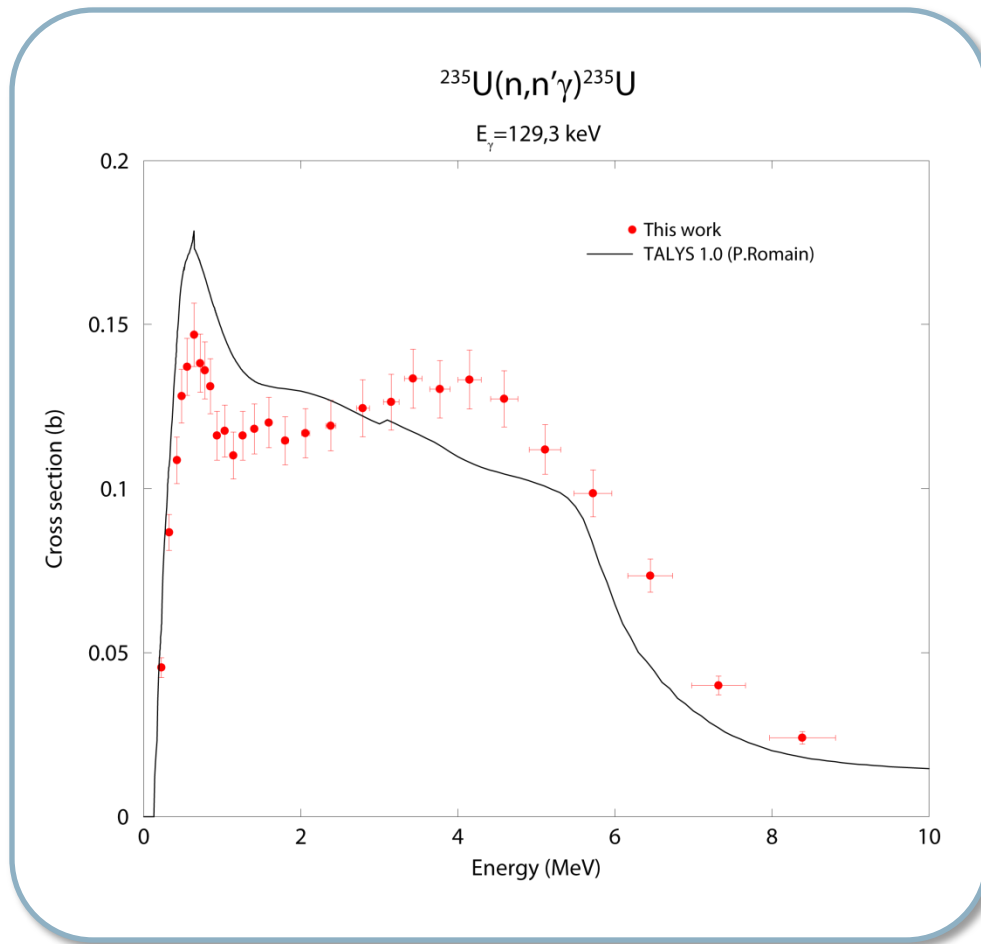
Results: $^{206}\text{Pb}(n,n'\gamma)^{206}\text{Pb}$



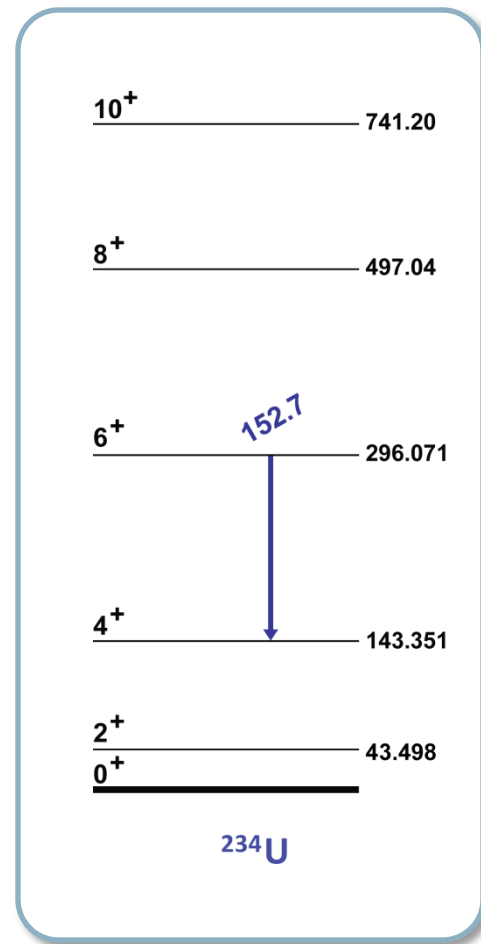
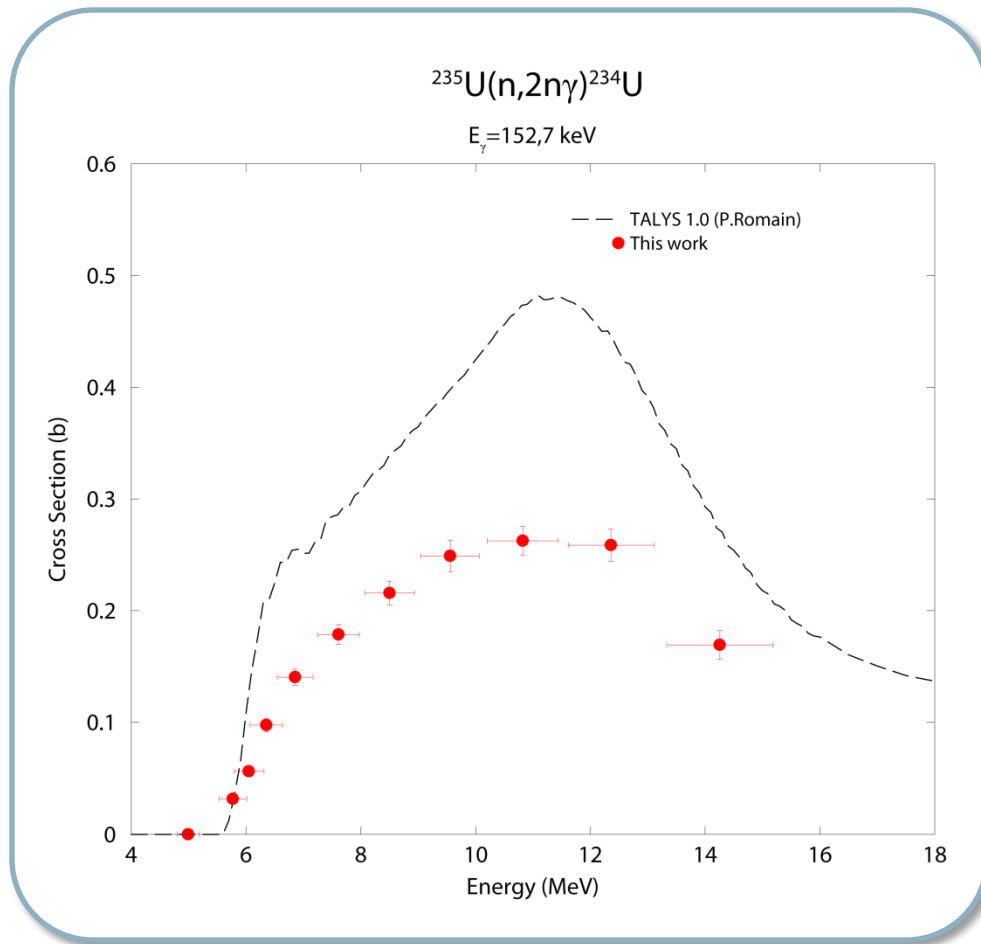
Results: $^{206}\text{Pb}(n,n'\gamma)^{206}\text{Pb}$



Results: $^{235}\text{U}(n,n'\gamma)^{235}\text{U}$



Results: $^{235}\text{U}(n,2n\gamma)^{234}\text{U}$



Conclusions & Perspectives

- Promising results for Pb and ^{235}U
- Further investigation on TALYS parameterization

- Perspectives:
 - Calibration measurements of the fission chamber
 - Development of a segmented planar HPGe detector, to reduce pile up rates
 - Measuring $(n, xn\gamma)$ reactions on actinides:
 - ^{232}Th : $(n, n'\gamma)$, $(n, 3n\gamma)$
 - ^{238}U : $(n, n'\gamma)$
 - ^{233}U : $(n, n'\gamma)$, $(n, 2n\gamma)$

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Angle integration

- One can show that the differential cross section can be expressed as a finite sum over even degree Legendre polynomials:

$$\frac{d\sigma}{d\Omega}(\theta_i, E_i) = \frac{\sigma}{4\pi} \sum_{k=0}^{[j_i]} c_{2k} P_{2k}(\cos \theta_i)$$

- where $[j_i]$ is the largest integer smaller than or equal to the spin of the decaying state j_i . Gaussian quadrature specialized to even degree polynomials [4] can be used: for a given number n of detectors, we need the best angles ($x_i = \cos \vartheta_j$) and weights w_i to verify:

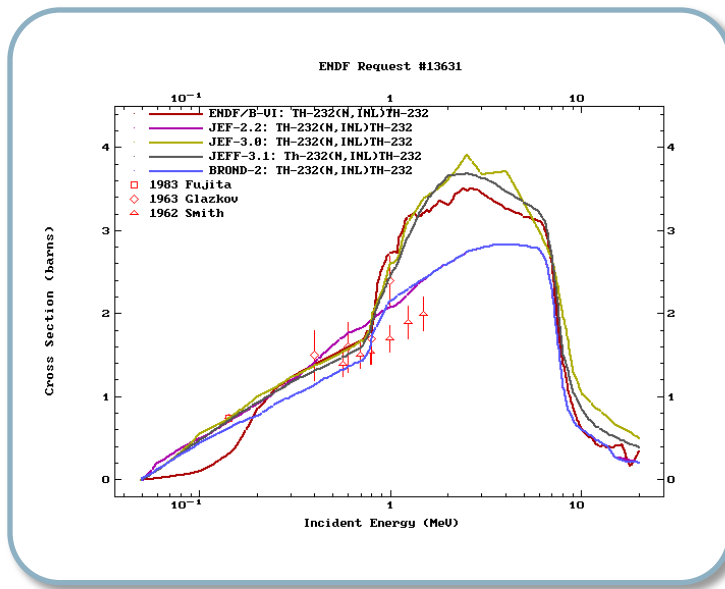
$$\sigma = 2\pi \int_{-1}^1 \frac{d\sigma}{d\Omega}(x) dx = 2\pi \sum_{i=1}^n w_i \frac{d\sigma}{d\Omega}(x_i)$$

- This is achieved for x_i being the n positive zeroes of the Legendre polynomial P_{2n} and

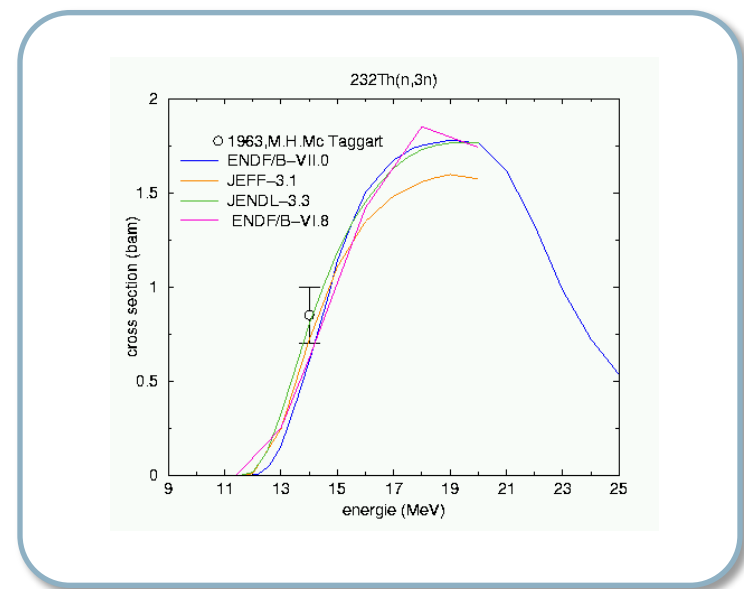
$$w_i = \frac{1}{\sum_{m=0}^{n-1} ((4m+1)/2) P_{2m}^2(x_i)}$$

- For the two detector configuration, the following results are found:
 $w = 0.69571$ for $\vartheta = 30.56^\circ, 149.44^\circ$ and $w = 1.30429$ for $\vartheta = 70.12^\circ, 109.88^\circ$

Currently known ^{232}Th data



$^{232}\text{Th}(n,\text{INL})$



$^{232}\text{Th}(n,3n)$